

Numerical methods of the fractional calculus

(continued II)

1

Computation of the Mittag-Leffler function

(1) Directly using the definition:

```
function y=mltlef(alpha,beta,z,N)
% Evaluation of the Mittag-Leffler function in two parameters:
% E_alpha,beta(z), using its series expansion.
%
% PARAMETERS:
% alpha - first index (scalar), beta - second index (scalar)
% z - row vector of values of the function argument
% N - number of terms in the power expansion (scalar)
% OPTIONAL, default N=100.

if nargin<4, N=100; end

m1=max(size(z));
m2=min(size(z));
if m2>1, z=z(1,:); end

k= repmat((1:N)',1,m1);
t= repmat(z,N,1);
a=t.*(k-1);
y= repmat(gamma(alpha*(1:N)'-1)+beta, 1, m1);
y=sum(t./a);
```

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Computation of binomial coefficients

The following coefficients appear frequently (e.g., in G1):

$$w_k^{(\alpha)} = (-1)^k \binom{\alpha}{k}, \quad k = 0, 1, 2, \dots,$$

It is possible to use the following recurrence:

$$w_0^{(\alpha)} = 1; \quad w_k^{(\alpha)} = \left(1 - \frac{\alpha+1}{k}\right) w_{k-1}^{(\alpha)}, \quad k = 1, 2, 3, \dots$$

Possible MATLAB code:

```
function y=bcrecur(a, n)
y=cumprod([1, 1 - ((a+1) ./ [1:n])]);
```

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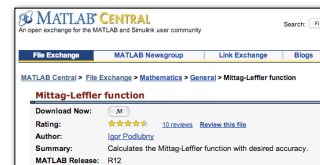
Computation of the Mittag-Leffler function

(2) There is a routine at MATLAB File Exchange:

<http://www.mathworks.com/matlabcentral/fileexchange/loadFile.do?objectId=87388&objectType=FILE>

MLF(alpha,beta,Z,P)

is the Mittag-Leffler function $E_{\alpha,\beta}(Z)$ evaluated with accuracy 10^{-P} for each element of Z.



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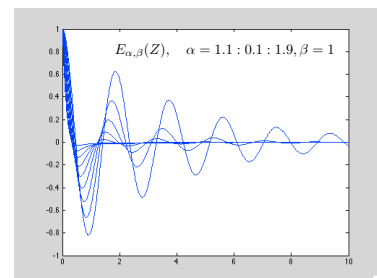
Computation of binomial coefficients

```
>> bcrecur(1.8,10)
ans =
Columns 1 through 8
1.0000 -1.8000 0.7200 0.0480 0.0144 0.0063 0.0034 0.0020
Columns 9 through 11
0.0013 0.0009 0.0007
>> bcrecur(2,10)
ans =
1 -2 1 0 0 0 0 0 0 0 0
>> bcrecur(3,10)
ans =
Columns 1 through 8
1.0000 -3.0000 3.0000 -1.0000 0 0 0 0
Columns 9 through 11
0 0 0
```

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Computation of the Mittag-Leffler function

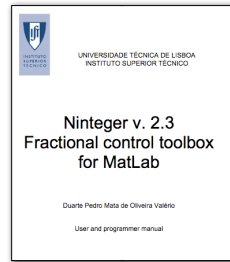
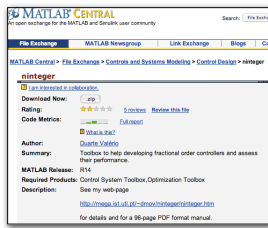
(2) There is a routine at MATLAB File Exchange:



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NINTEGER Toolbox by Duarte Valerio



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Sample applications

Fractional derivatives of order 1/2 can appear even within the context of the integer-order models!

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NINTEGER Toolbox by Duarte Valerio

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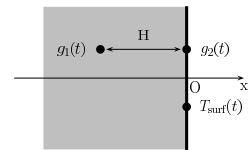
Heat flux in a blast furnace wall

$$\hat{c}\rho \frac{\partial T}{\partial t} = \hat{\lambda} \frac{\partial^2 T}{\partial x^2}, \quad (t > 0, \quad -\infty < x < 0)$$

$$T(0, x) = T_0$$

$$T(t, 0) = T_{\text{surf}}(t)$$

$$\left| \lim_{x \rightarrow -\infty} T(t, x) \right| < \infty$$



t - is time [s],

x - is the spatial coordinate in the direction of heat conduction [m],

\hat{c} - is heat capacity [$J \text{ kg}^{-1} \text{ K}^{-1}$],

$\hat{\rho}$ - is mass density [kg m^{-3}],

$T(t, x)$ - is temperature [K],

$\hat{\lambda}$ - is the coefficient of heat conductivity [$\text{W m}^{-1} \text{ K}^{-1}$]

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NINTEGER Toolbox by Duarte Valerio

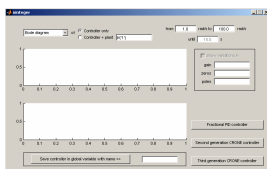


Figure 2. The main dialog of NinInteger before being filled in

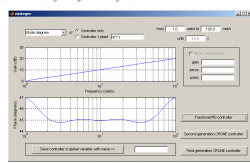


Figure 3. The main dialog displaying a block diagram

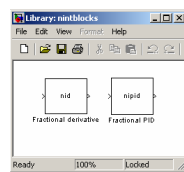


Figure 4. Simulink library

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Heat flux in a blast furnace wall

Introduce an auxiliary function

$$u(t, x) = T(t, x) - T_0.$$

which must be a solution of the following problem:

$$\hat{c}\rho \frac{\partial u}{\partial t} = \hat{\lambda} \frac{\partial^2 u}{\partial x^2}, \quad (t > 0, \quad -\infty < x < 0)$$

$$u(0, x) = 0$$

$$u(t, 0) = T_{\text{surf}}(t) - T_0$$

$$\left| \lim_{x \rightarrow -\infty} u(t, x) \right| < \infty$$

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Heat flux in a blast furnace wall

The Laplace transform gives:

$$s \hat{c} \hat{\rho} U(s, x) = \hat{\lambda} \frac{d^2 U(s, x)}{dx^2}$$

The solution bounded at $x \rightarrow -\infty$ is:

$$U(s, x) = U(s, 0) \exp \left(x \sqrt{\frac{\hat{c} \hat{\rho} s}{\hat{\lambda}}} \right)$$

Then we have

$$\frac{dU}{dx}(s, x) = U(s, 0) \sqrt{\frac{\hat{c} \hat{\rho} s}{\hat{\lambda}}} \exp \left(x \sqrt{\frac{\hat{c} \hat{\rho} s}{\hat{\lambda}}} \right)$$

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Heat flux in a blast furnace wall

Numerical solution: short memory principle

$$q_A(t) \approx \tilde{q}_A(t) = \sqrt{\hat{c} \hat{\rho} \hat{\lambda}} {}_{(t-L)} D_t^{1/2} g(t)$$

Normalized error:

$$\delta_0 = \frac{|q_A(t) - \tilde{q}_A(t)|}{M} = \frac{1}{\sqrt{L} \Gamma(\frac{1}{2})}, \quad M = \max_{[0, \infty]} |g(t)|$$

Condition on the memory length: $L \geq \frac{1}{\pi \delta_0^2}$

Calculate using:

$${}_{(t-L)} D_t^{1/2} g(t) = \tau^{-\alpha} \sum_{i=0}^{N(t)} c_i g(t - i\tau),$$

$$N(t) = \min \left\{ \left\lceil \frac{t}{\tau} \right\rceil, \left\lceil \frac{L}{\tau} \right\rceil \right\}, \quad c_i = (-1)^i \binom{1/2}{i}$$

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Heat flux in a blast furnace wall

$$U(s, x) = U(s, 0) \exp \left(x \sqrt{\frac{\hat{c} \hat{\rho} s}{\hat{\lambda}}} \right) \quad \frac{dU}{dx}(s, x) = U(s, 0) \sqrt{\frac{\hat{c} \hat{\rho} s}{\hat{\lambda}}} \exp \left(x \sqrt{\frac{\hat{c} \hat{\rho} s}{\hat{\lambda}}} \right)$$

$$\frac{1}{\sqrt{s}} \frac{dU}{dx}(s, 0) = \sqrt{\frac{\hat{c} \hat{\rho}}{\hat{\lambda}}} U(s, 0)$$

The inverse Laplace transform gives:

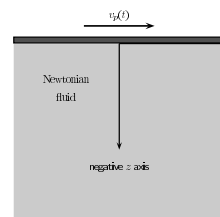
$${}_0 D_t^{-1/2} \frac{\partial u}{\partial x}(t, 0) = \sqrt{\frac{\hat{c} \hat{\rho}}{\hat{\lambda}}} u(t, 0)$$

After fractional integration of both sides:

$$\frac{\partial u}{\partial x}(t, 0) = \sqrt{\frac{\hat{c} \hat{\rho}}{\hat{\lambda}}} {}_0 D_t^{1/2} u(t, 0)$$

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A plate on a Newtonian fluid



The equation of motion of the fluid is

$$\rho \frac{\partial v}{\partial t} = \mu \frac{\partial^2 v}{\partial z^2}, \quad (0 < t < \infty, \quad -\infty < z < 0)$$

assume that initially the fluid is in equilibrium state:

$$v(0, z) = 0, \quad (-\infty < z < 0)$$

the influence of the plate's motion vanishes for $z \rightarrow \infty$:

$$v(t, -\infty) = 0, \quad (0 < t < \infty)$$

The fluid's velocity at $z = 0$ is equal to the given velocity of the plate:

$$v(t, 0) = v_p(t)$$

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Heat flux in a blast furnace wall

$$\frac{\partial u}{\partial x}(t, 0) = \sqrt{\frac{\hat{c} \hat{\rho}}{\hat{\lambda}}} {}_0 D_t^{1/2} u(t, 0) \quad u(t, x) = T(t, x) - T_0,$$

$$q_A(t) = \sqrt{\hat{c} \hat{\rho} \hat{\lambda}} {}_0 D_t^{1/2} g(t), \quad g(t) = T_{\text{surf}}(t) - T_0$$

$$q_A(t) = \hat{\lambda} \frac{\partial T}{\partial x}(t, 0)$$

the heat flux

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A plate on a Newtonian fluid

Applying the Laplace transform gives:

$$\rho s V(s, z) = \mu \frac{d^2 V(s, z)}{dz^2}$$

$$V(s, 0) = V_p(s),$$

$$V(s, -\infty) = 0,$$

The solution of this problem is:

$$V(s, z) = V_p(s) \exp \left(z \sqrt{\frac{\rho s}{\mu}} \right)$$

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A plate on a Newtonian fluid

By differentiation we obtain:

$$\frac{dV(s,z)}{dz} = \sqrt{\frac{\rho s}{\mu}} V_p(s) \exp\left(z \sqrt{\frac{\rho s}{\mu}}\right) = \sqrt{\frac{\rho s}{\mu}} V(s,z)$$

Knowing the velocity, we can obtain the shear stress:

$$\sigma(t,z) = \mu \frac{\partial v(t,z)}{\partial z}$$

which in terms of the Laplace transforms reads:

$$\bar{\sigma}(s,z) = \mu \frac{dV(s,z)}{dz} = \sqrt{\mu \rho s} V(s,z).$$

and therefore, in time domain we have:

$$\sigma(t,z) = \sqrt{\mu \rho} D_t^{1/2} v(s,z)$$

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Motion of an immersed plate

Sample numerical solution

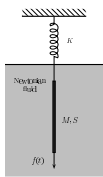
```
% (2) Make the main matrix B2_{N}^{(2)}:
alpha=2;
bc=flipr(bcrecur(alpha,N-1));
for k=1:N
    B2(k,1:k)=bc((N-k+1):N);
end

% (3) Make the matrix for the entire equation:
B=h^(-2)*B2 + h^(-1.5)*B1 + eye(size(B1));

% Make right-hand side:
F=8*(Tc=1)';
```

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Motion of an immersed plate



Considering the forces we have:

$$My''(t) = f(t) - Ky(t) - 2S\sigma(t,0)$$

Using the relationships

$$\sigma(t,z) = \sqrt{\mu \rho} D_t^{1/2} v(s,z) \quad v_p(t,0) = y'(t)$$

we arrive at the following equation:

$$Ay''(t) + B_0 D_t^{3/2} y(t) + Cy(t) = f(t) \quad (t > 0),$$

$$A = M, \quad B = 2S\sqrt{\mu \rho}, \quad C = K,$$

add initial conditions! $y(0) = 0, \quad y'(0) = 0$

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Motion of an immersed plate

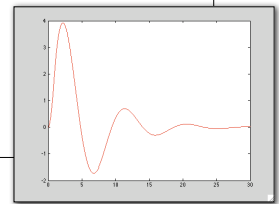
Sample numerical solution

```
% Make right-hand side:
F=8*(Tc=1)';

% (3) Use zero initial conditions:
F(1)=0;
F(2)=0;
B(1,1)=1;
B(2,1)=-1/h; B(2,2)=1/h;

% (4) Solve the system BY=F:
Y=B\F;

% Plot the solution:
figure(1)
plot(T,Y,'r')
```



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Motion of an immersed plate

Sample numerical solution

```
% Script for solving Bagley-Torvik equation with zero
initial conditions:

clear all

alpha=1.5;

% Numerical solution:
h=0.075;
T=0:h:30;
N=30/h + 1;
B=zeros(N,N);

% (1) Make the main matrix B1_{N}^{(alpha)}:
bc=flipr(bcrecur(alpha,N-1));
for k=1:N
    B1(k,1:k)=bc((N-k+1):N);
end
```

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Short memory principle for FDEs

Consider the problem:

$${}_0 D_t^{3/2} y(t) + y(t) = f(t), \quad (t > 0)$$

$$y(0) = y'(0) = 0$$

for the following RHS:

1. $f(t) \equiv 1$
2. $f(t) = te^{-t}$
3. $f(t) = t^{-1}e^{-1/t}$
4. $f(t) = e^{-t} \sin(0.2t)$

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Short memory principle for FDEs

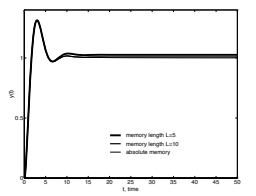


Figure 8.8: Solution of the problem (8.56) for $f(t) \equiv 1$

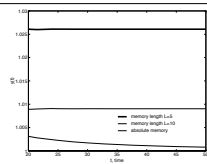


Figure 8.9: Solution of the problem (8.56) for $f(t) \equiv 1$ (zoom)

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Short memory principle for FDEs

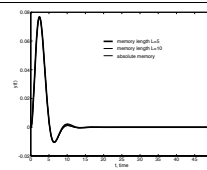


Figure 8.11: Solution of the problem (8.56) for $f(t) = e^{-t} \sin(0.2t)$

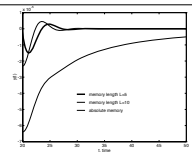


Figure 8.12: Solution of (8.56) for $f(t) = e^{-t} \sin(0.2t)$ (zoom)

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Short memory principle for FDEs

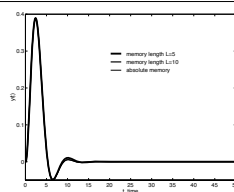


Figure 8.10: Solution of the problem (8.56) for $f(t) = te^{-t}$

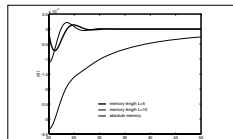


Figure 8.11: Solution of the problem (8.56) for $f(t) = te^{-t}$ (zoom)

Short memory principle for FDEs

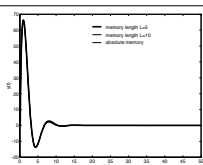


Figure 8.12: Solution of the problem (8.56) for $f(t) = t^{-1/2} e^{-t/2}$

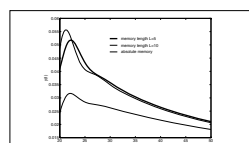


Figure 8.13: Solution of the problem (8.56) for $f(t) = t^{-1/2} e^{-t/2}$ (zoom)

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